

1. Mass renormalization in ϕ^4 theory

Let's study renormalization of the mass parameter in ϕ^4 theory. As in the notes we consider

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 - \frac{1}{4!} \lambda \phi^4$$

with a cutoff on the Euclidean loop momentum Λ , and

$$\mathcal{L}' = \frac{1}{2} \partial_\mu \phi' \partial^\mu \phi' - \frac{1}{2} m'^2 \phi'^2 - \frac{1}{4!} \lambda' \phi'^4$$

with a cutoff $\Lambda' = \Lambda - \delta\Lambda$. The idea is to match the tree-level propagator in the primed theory to the corresponding quantity in the unprimed theory, namely the sum of diagrams



(i) In the primed theory the propagator is

$$\frac{i}{p^2 - m'^2}.$$

Set $m'^2 = m^2 + \delta m^2$ where $\delta m^2 = -\frac{dm^2}{d\Lambda} \delta\Lambda$. Expand the propagator to first order in $\delta\Lambda$.

(ii) Match your answer to the corresponding calculation in the unprimed theory. You can stop at a single χ loop, and for simplicity you can assume $m \ll \Lambda$. You should obtain a trivial differential equation for the mass parameter and solve it to find $m^2(\Lambda)$.

2. Pion decay

- (i) Derive the Noether currents $j_L^{\mu a}, j_R^{\mu a}$ associated with the $SU(2)_L \times SU(2)_R$ symmetry

$$\delta\psi_L = -\frac{i}{2}\lambda_L^a \sigma^a \psi_L \quad \delta\psi_R = -\frac{i}{2}\lambda_R^a \sigma^a \psi_R$$

of QCD with two flavors of massless quarks. We'll mostly be interested in the vector and axial-vector linear combinations $j_V^{\mu a} = j_L^{\mu a} + j_R^{\mu a}$, $j_A^{\mu a} = -j_L^{\mu a} + j_R^{\mu a}$.

- (ii) Repeat part (i) for the $SU(2)$ non-linear σ -model

$$\mathcal{L} = \frac{1}{4}f^2 \text{Tr} (\partial_\mu U^\dagger \partial^\mu U)$$

where the symmetry is

$$\delta U = -\frac{i}{2}\lambda_L^a \sigma^a U + U \frac{i}{2}\lambda_R^a \sigma^a.$$

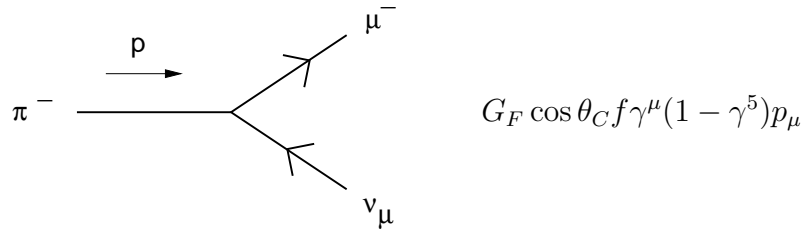
You only need to work out the symmetry currents to first order in the pion fields π^a , where $U = e^{i\pi^a \sigma^a / f}$.

- (iii) The weak interaction responsible for the decay $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$ is

$$\mathcal{L}_{\text{weak}} = -\frac{1}{\sqrt{2}}G_F \cos \theta_C \bar{\mu} \gamma^\lambda (1 - \gamma^5) \nu_\mu \bar{u} \gamma_\lambda (1 - \gamma^5) d + h.c.$$

Here $\theta_C \approx 13^\circ$ is the ‘Cabibbo angle.’ Suppose we can identify the symmetry currents worked out in parts (i) and (ii). Use this to rewrite $\mathcal{L}_{\text{weak}}$ in terms of the fields μ, ν_μ, π^a , again working to first order in the pion fields.

- (iv) If I did it right this leads to a vertex



Calculate the pion lifetime in terms of G_F , θ_C , f , m_π , m_μ . Use the observed lifetime $\tau_\pi = 2.6 \times 10^{-8}$ sec to estimate f .

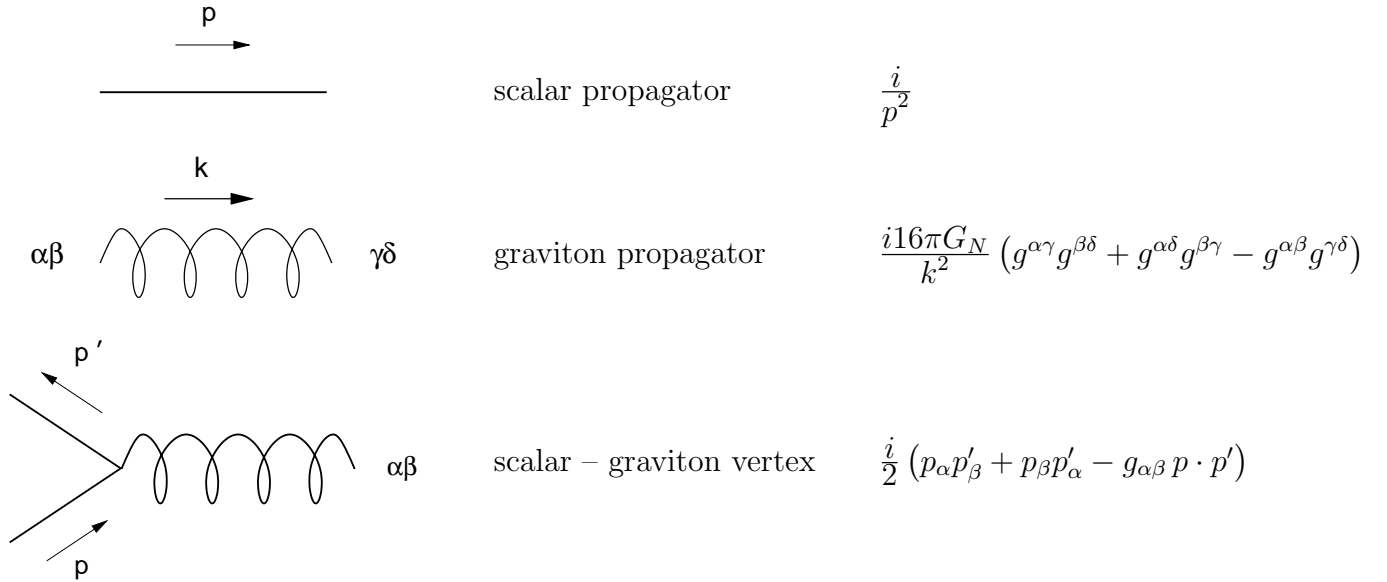
- (v) The decay $\pi^- \rightarrow e^- \bar{\nu}_e$ only differs by replacing $\mu \rightarrow e$, $\nu_\mu \rightarrow \nu_e$. Predict the branching ratio

$$\frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu)}$$

How well did you do?

3. Unitarity violation in quantum gravity

Consider two distinct types of massless scalar particles A and B which only interact gravitationally. The Feynman rules are



Here G_N is Newton's constant and $g_{\alpha\beta} = \text{diag}(+---)$ is the Minkowski metric. The vertices and propagators are the same for A and B .

- (i) Compute the tree-level amplitude and center-of-mass differential cross section for the process $AA \rightarrow BB$.

- (ii) The partial-wave expansion of the scattering amplitude is

$$f(\theta) = \frac{1}{i\sqrt{s}} \sum_{l=0}^{\infty} (2l+1) P_l(\cos \theta) S_l(E)$$

where P_l is a Legendre polynomial. This is related to the center-of-mass differential cross section by

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{cm}} = |f(\theta)|^2.$$

Compute the partial-wave S-matrix elements $S_l(E)$. For which values of l are they non-zero?

- (iii) At what center-of-mass energy is the unitarity bound $|S_l(E)| \leq 1$ violated?